

# (1-x)Li<sub>4</sub>WO<sub>5</sub>-xLiF: A novel oxyfluoride system and their microwave dielectric properties

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## ARTICLE INFO

### Article history:

Received 9 December 2019

Received in revised form

21 April 2020

Accepted 21 April 2020

Available online 28 April 2020

### Keywords:

Degree of order

Relative density

Microwave dielectric properties

Rock salt structure

## ABSTRACT

A series of novel oxyfluorides in the (1-x)Li<sub>4</sub>WO<sub>5</sub>-xLiF ( $x = 0.1, 0.2, 0.3, 0.4, 0.5$ ) binary system were prepared by the conventional solid-state ceramic route and their microwave dielectric properties were reported. Phase transformation from cubic to triclinic was confirmed via XRD, DTA and Raman spectra characterization. The temperature of phase transition (cubic-triclinic phase) moved towards high temperature with the increase of LiF content. EDS analysis and element mapping implied the homogeneous distribution of F<sup>-</sup> ions. Particularly, when  $x$  was equal to 0.3 (The chemical formula was Li<sub>3.1</sub>W<sub>0.7</sub>F<sub>0.3</sub>O<sub>3.5</sub>), the microstructure of Li<sub>4</sub>WO<sub>5</sub> was evidently improved and a dense microstructure with high relative density (96.4%) was achieved. The comprehensive microwave dielectric performance of Li<sub>3.1</sub>W<sub>0.7</sub>F<sub>0.3</sub>O<sub>3.5</sub> ceramics sintered at 930 °C was optimum:  $Q \times f = 71,000$  GHz,  $\epsilon_r = 16.6$  and  $\tau_f = -16.4$  ppm/°C. The infrared reflectivity spectra revealed that the intrinsic quality factor of Li<sub>3.1</sub>W<sub>0.7</sub>F<sub>0.3</sub>O<sub>3.5</sub> ceramics was 83,000 GHz.

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## 1. Introduction

Recently, the dielectric properties of Li-containing rock salt compounds have received considerable attention because of their ultralow dielectric loss at microwave frequency region [1–8]. Unfortunately, the sintering temperature (>960 °C) of most rock salt compounds is too high to be directly used in low-temperature cofired ceramic technology (LTCC), such as Li<sub>2</sub>MnO<sub>3</sub> and Li<sub>2</sub>Mg-TiO<sub>4</sub>. So it is necessary to add sintering aids to decrease the sintering temperature [1,3]. The addition of sintering aids not only reduces the performance and complicate process, but also increases the possibility of reacting with the electrode [9–11]. Among them, Li<sub>4</sub>WO<sub>5</sub> with the sintering temperature of 890 °C was reported to be

a low-firing ceramic with promising microwave dielectric properties [8]. However, it is still a challenge to obtain well-sintered Li<sub>4</sub>WO<sub>5</sub> ceramics with high density and low porosity for optimizing the dielectric properties, which relates to the phase transition and the Li volatilization during the high-temperature sintering process. In our previous work, non-stoichiometry with addition of extra Li were carried out in Li<sub>4</sub>WO<sub>5</sub> to suppress Li volatilization, which not only improved the densification of Li<sub>4</sub>WO<sub>5</sub> but also enhanced the dielectric performances [12,13]. It is concluded that lowering the sintering temperature is crucial to reduce the possibility of chemical elements volatilization. To date, the influence of phase transition on microstructural characteristics and macroscopic physical properties has not yet been reported.

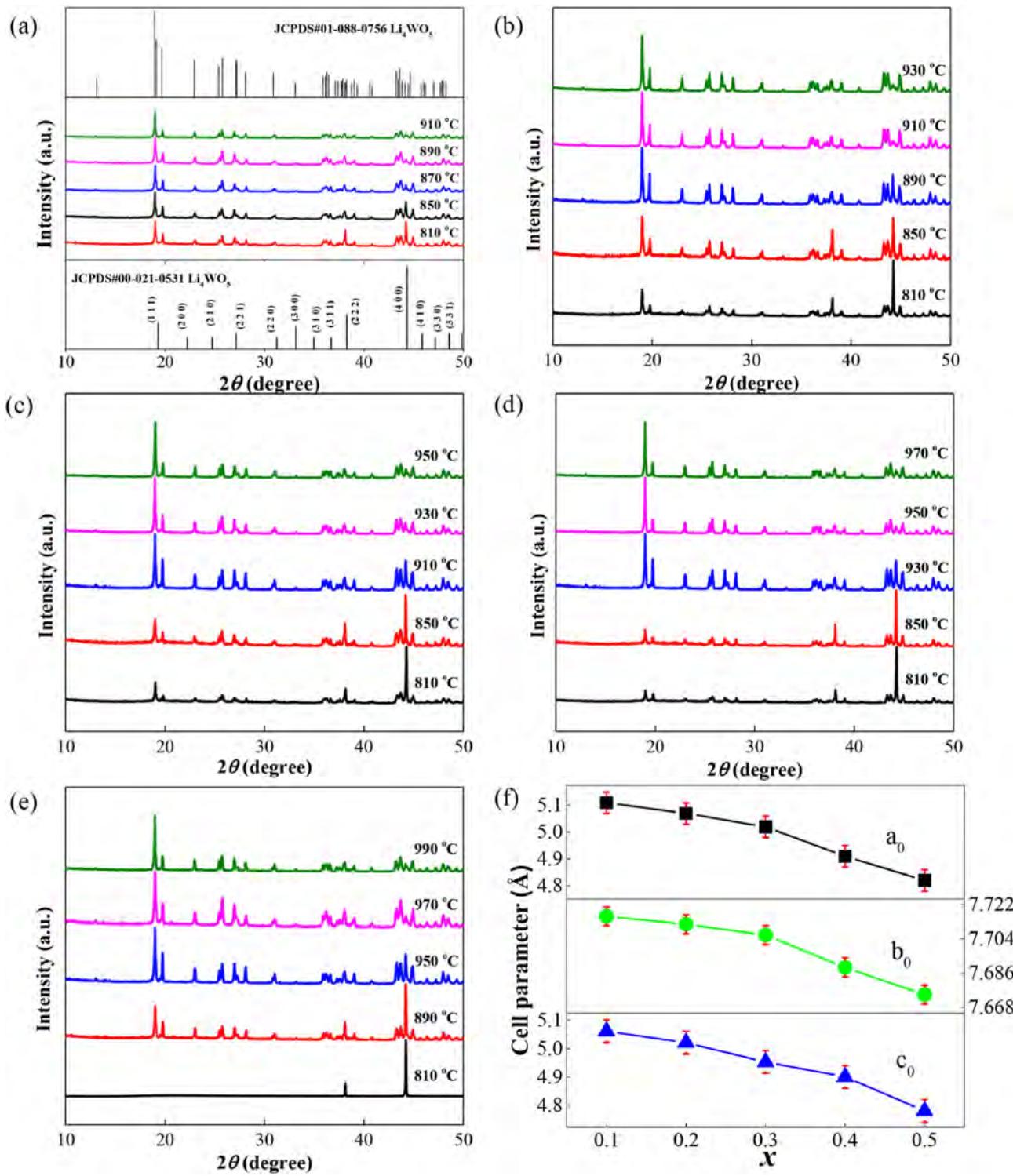
Addition of sintering aids (B<sub>2</sub>O<sub>3</sub>, MnO<sub>2</sub> and LiF) with low melting points is a promising and effective approach to lower the sintering temperatures of ceramics [14–16]. Herein, LiF was selected as sintering aid for Li<sub>4</sub>WO<sub>5</sub> due to its low melting point (848 °C) and successful reduction of sintering temperatures of some materials, such as Li<sub>2</sub>TiO<sub>3</sub>, Li<sub>3</sub>NbO<sub>4</sub> [17,18]. In addition, LiF adopts the same rock salt structure with Li<sub>4</sub>WO<sub>5</sub>, making it possible to form solid solution. Castellanos et al. reported that Li<sub>2</sub>TiO<sub>3</sub> could form solid solutions with MgO and the phase transition temperature from

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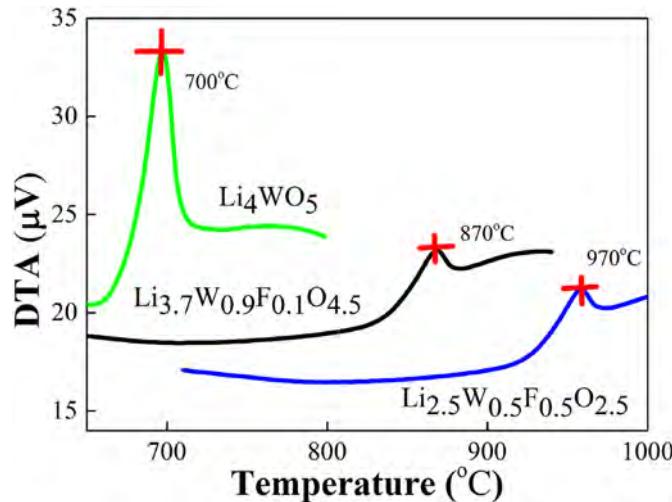
**Fig. 1.** (a–e) The room-temperature XRD patterns of the  $(1-x)\text{Li}_4\text{WO}_5-x\text{LiF}$  ((a)  $x = 0.1$ , (b)  $x = 0.2$ , (c)  $x = 0.3$ , (d)  $x = 0.4$ , (e)  $x = 0.5$ ) ceramics sintered from 810 °C to 990 °C for 4 h; Fig. 1(f) the cell parameters of the triclinic phase as a function of LiF content sintered at optimum temperature.

monoclinic to cubic increased rapidly with the increase of MgO content [19]. Similarly,  $\text{Li}_4\text{WO}_5$  might form solid solution with LiF and the phase transition temperature from cubic to triclinic might be changed. If the addition of LiF could decrease the degree of order and the sintering temperature of  $\text{Li}_4\text{WO}_5$ , the cubic phase solid solutions might be obtained. Therefore, new oxyfluorides in rock salt family are expected to obtain. In this work, a series of oxyfluorides in the  $(1-x)\text{Li}_4\text{WO}_5-x\text{LiF}$  ( $x = 0.1, 0.2, 0.3, 0.4, 0.5$ ) binary

system were prepared. The phase evolution, sintering behavior, and microwave dielectric properties were studied in detail.

## 2. Experimental

The starting materials of  $\text{Li}_2\text{CO}_3$  (99.99%),  $\text{WO}_3$  (99%) were used according to the stoichiometric formula  $\text{Li}_4\text{WO}_5$ . The raw materials were ball milled for 6 h with ethanol in a nylon jar, followed by a



**Fig. 2.** DTA curves for  $\text{Li}_4\text{WO}_5$ ,  $\text{Li}_{3.7}\text{W}_{0.9}\text{F}_{0.1}\text{O}_{4.5}$  and  $\text{Li}_{2.5}\text{W}_{0.5}\text{F}_{0.5}\text{O}_{2.5}$  powders.

treatment of drying and calcination at  $600^\circ\text{C}$  for 4 h. Powders of  $(1-x)\text{Li}_4\text{WO}_5-x\text{LiF}$  ( $x = 0.1, 0.2, 0.3, 0.4, 0.5$ ) were ball milled in alcohol for 6 h, then, dried and the obtained powders were mixed with 5 wt % of polyvinyl alcohol (PVA) as a binder and granulated. The resultant powders were pressed into the disks with 12 mm diameter and 6 mm thickness at 200 MPa. Pellets were sintered in air at temperature ranges from  $800^\circ\text{C}$  to  $1100^\circ\text{C}$  at a heating rate of  $5^\circ\text{C}/\text{min}$ .

The crystal structure and phase purity of the sintered samples were identified using X-ray diffraction with  $\text{CuK}_\alpha$  radiation (XRD; Model X'Pert PRO, PANalytical, Almelo, the Netherlands) at a scanning rate of  $0.02^\circ\text{s}^{-1}$  in a  $2\theta$  range of  $10^\circ\text{--}50^\circ$ . The ratio of cubic/triclinic phase of  $\text{Li}_4\text{WO}_5-x\text{LiF}$  was calculated by the original XRD data. And the ratio of the main peak of cubic phase and triclinic phase represents the ratio of cubic and triclinic phase, respectively. The phase composition of the sintered ceramics was also confirmed by Raman spectrometer (DXR; Thermo Fisher Scientific, American). Selected area electron diffraction (SAED) data was performed with

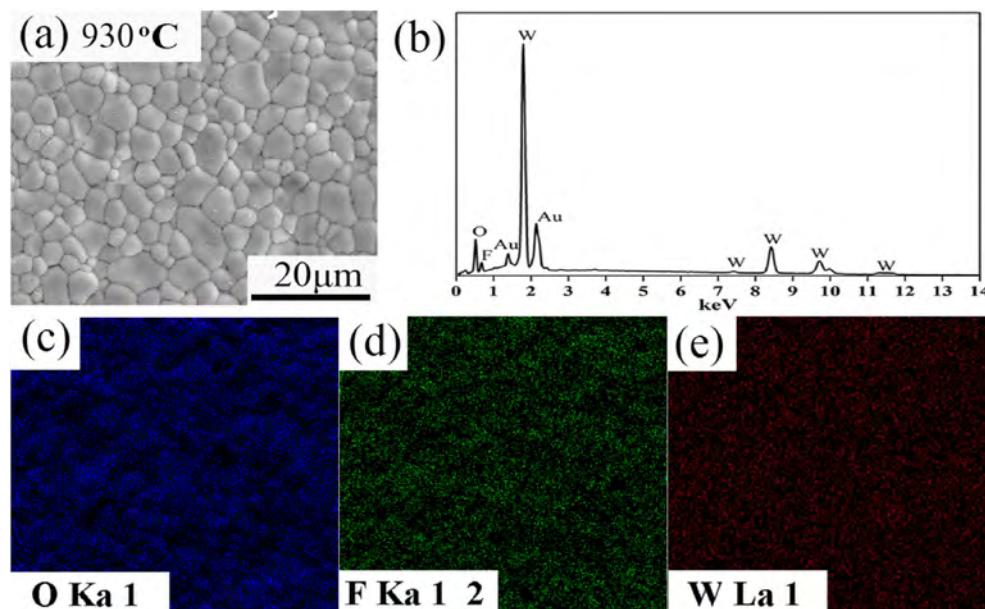
a JEM-2100 F at an accelerating voltage of 200 kV. The bulk densities were measured by Archimedes' method. To observe grain morphology, the polished and thermally etched surfaces of fired ceramics (fired  $25^\circ\text{C}$  below the optimized sintering temperature) were observed with a Hitachi (Tokyo, Japan) S4800 field-emission scanning electron microscopy (FESEM). Meanwhile, the elemental mapping and energy-dispersive x-ray spectroscopy (EDS) were also measured by FESEM. The permittivity and the quality factor values were measured in the TE<sub>011</sub> mode according to the Hakki-Coleman method using an Agilent N5230A network analyzer (Palo Alto, CA) in the temperature range of  $25^\circ\text{C}$ – $85^\circ\text{C}$  [20]. The temperature coefficient of the resonant frequency  $\tau_f$  was calculated by Equation (1):

$$\tau_f = \frac{\Delta f}{f_1(T_2 - T_1)} \quad (1)$$

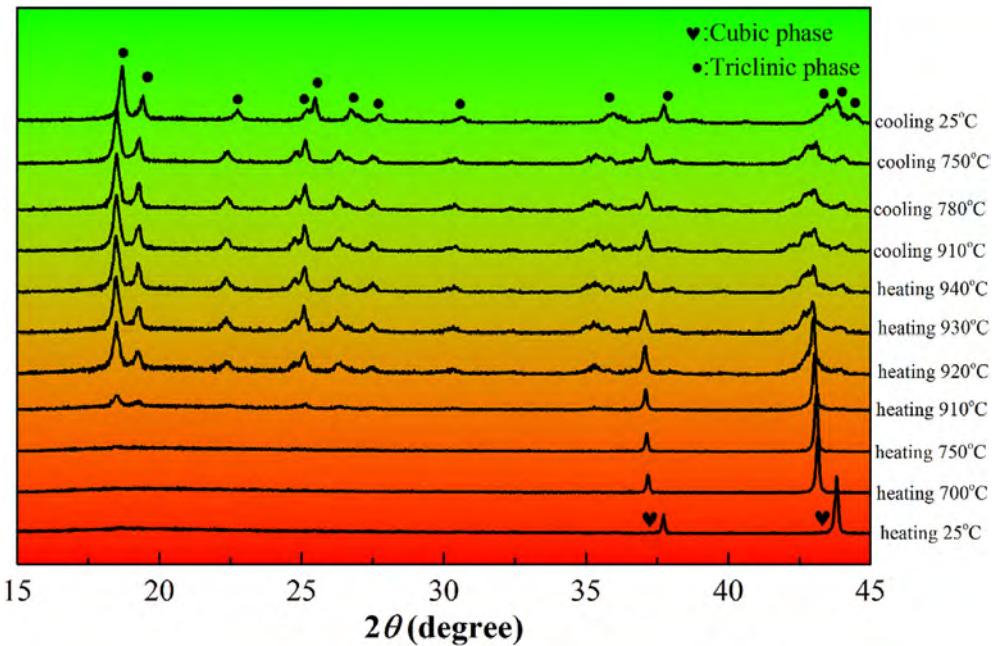
where  $\Delta f = f_2 - f_1$ ,  $f_2$  and  $f_1$  were the resonant frequencies at the temperature of  $T_2$  ( $85^\circ\text{C}$ ) and  $T_1$  ( $25^\circ\text{C}$ ), respectively. The linear coefficient of thermal expansion (CTE) of the  $\text{Li}_4\text{WO}_5$  and  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramics were estimated utilizing a thermal dilatometer (DIL402C, NETZSCH, Germany). The infrared spectra at room temperature were measured by a Bruker IFS 66v FT-IR spectrometer on the infrared beamline station (U4) at the National Synchrotron Radiation Lab (NSRL), China.

### 3. Result and discussions

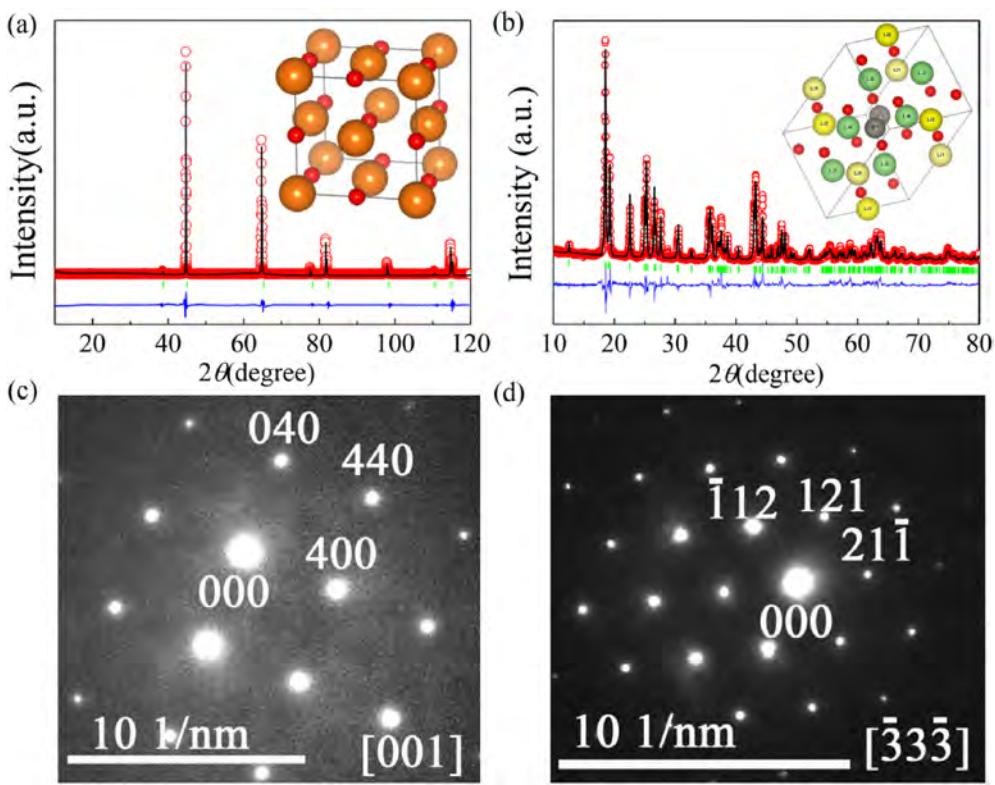
**Fig. 1(a–e)** shows the room-temperature XRD patterns of the  $(1-x)\text{Li}_4\text{WO}_5-x\text{LiF}$  ((a)  $x = 0.1$ , (b)  $x = 0.2$ , (c)  $x = 0.3$ , (d)  $x = 0.4$ , (e)  $x = 0.5$ ) ceramics sintered from  $810^\circ\text{C}$  to  $990^\circ\text{C}$  for 4 h. According to the previous reports, the  $\text{Li}_4\text{WO}_5$  exhibited an irreversible phase transition from cubic to triclinic around  $700^\circ\text{C}$  [8]. In the present work, this phase transformation was also observed. And the temperature of phase transformation changed with the  $x$  value in **Fig. 1**. Both the cubic  $\text{Li}_4\text{WO}_5$  (JCPDS card No. 00-021-0531) and triclinic  $\text{Li}_4\text{WO}_5$  (JCPDS card No. 01-088-0756) phase coexisted with the cubic phase prominent when samples sintered at low temperature, e.g.  $850^\circ\text{C}$ . With the increase of sintering temperature, the fraction



**Fig. 3.** SEM image (a), EDS spectrum (b) and EDS elemental mapping images (c) (d) (e) of the surface of  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramic sintered at  $930^\circ\text{C}$  for 4 h.



**Fig. 4.** The various-temperature XRD analysis was conducted on the  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramics sintered from room temperature to 950 °C.

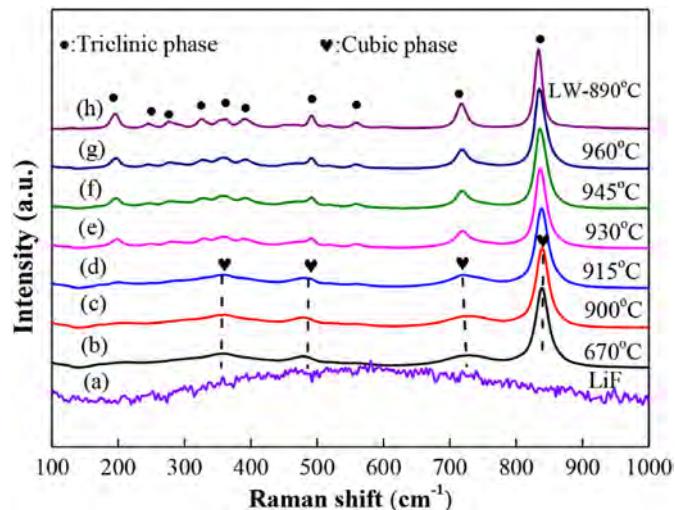


**Fig. 5.** (a) and (b) Rietveld refinement patterns and schematics of crystal structure for the cubic phase and triclinic phase of  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$ , respectively; (c) and (d) selected area electron diffraction (SAED) images for the cubic phase and triclinic phase of  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$ , respectively.

of the cubic phase gradually decreased. The cubic phase could fully transform to the triclinic phase ultimately at a critical temperature which depended on the LiF content. Fig. S1 shows the evolution in the ratio of cubic/triclinic phases (the dashed lines for the triclinic  $\gamma$  phase, and the solid lines for the cubic  $\alpha$  phase) with respect to  $x$  value. The DTA curves show the phase transformation temperature

increases with the increase of LiF content in Fig. 2, such as 700 °C for un-doped sample, 870 °C for  $\text{Li}_{3.7}\text{W}_{0.9}\text{F}_{0.1}\text{O}_{4.5}$  ( $x = 0.1$ ) ceramics and 970 °C for  $\text{Li}_{2.5}\text{W}_{0.5}\text{F}_{0.5}\text{O}_{2.5}$  ( $x = 0.5$ ) ceramics.

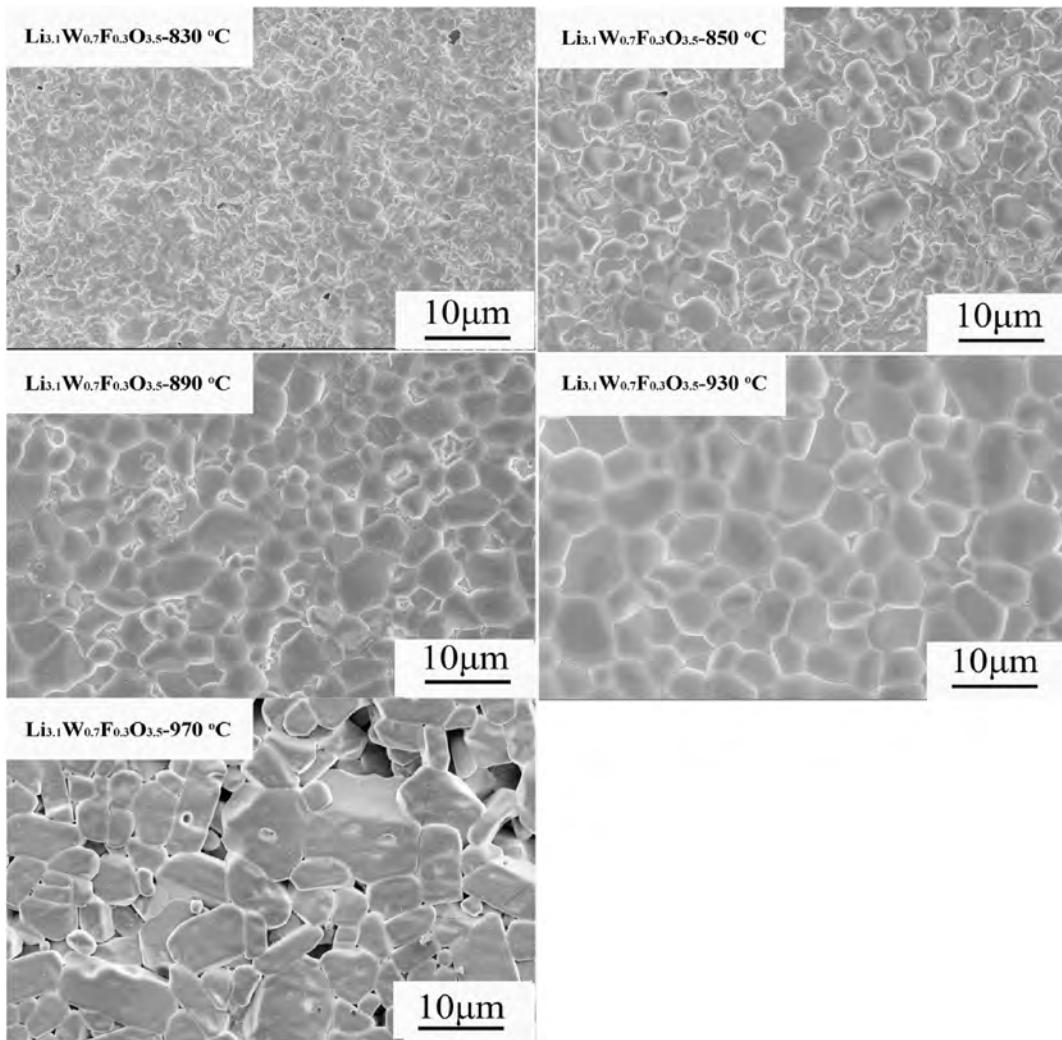
In the present system, two possible accommodation mechanisms of addition of LiF in the ceramic matrix are considered: (1) LiF might aggregate around the grain boundaries and facilitate



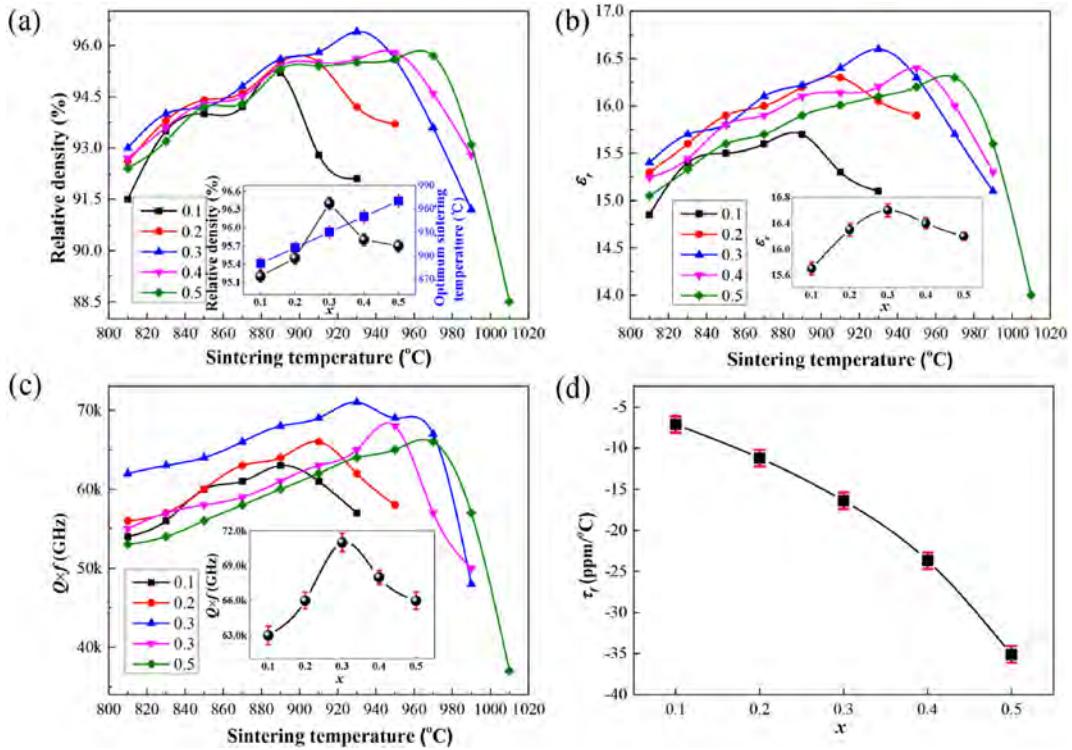
**Fig. 6.** The Raman spectra of LiF,  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  sintered from 670 °C to 960 °C and pure  $\text{Li}_4\text{WO}_5$  sintered at 890 °C, respectively.

densification through accelerating mass transfer process as sintering aids [16,17]. (2) LiF incorporated into the crystal lattice of  $\text{Li}_4\text{WO}_5$  to form solid solutions; or both. In the view of compositional dependence of phase transformation temperature, LiF enters into  $\text{Li}_4\text{WO}_5$  inducing evident effects on the degree of order of crystal structure. Based on the empirical Vegard's law, the lattice parameter of a solid solution is approximately linearly related to the molar fraction. Fig. 1(f) presents the variation in cell parameters of the triclinic phase as a function of LiF content. A nearly linear change is observed, providing direct evidence to verify the fact that LiF enters into the lattice of  $\text{Li}_4\text{WO}_5$ . The declining trend is attributed to the smaller ionic radius of  $\text{F}^-$  (1.19 Å) with respect to  $\text{O}^{2-}$  (1.26 Å). The SEM image of  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramics sintered at 930 °C shows a homogeneous and dense grain distribution (Fig. 3(a)). The element mapping analysis of  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramics is displayed in Fig. 3(c)–(e). The F element was distributed homogeneously in the entire region, and no element enrichment was observed at the grain boundary, indicating that LiF diffused into the lattice of  $\text{Li}_4\text{WO}_5$ . The results were consistent with those of the XRD and verified the solid solution formation.

In order to deeply understand the influence of LiF doping on phase transition of  $\text{Li}_4\text{WO}_5$ , various-temperature XRD analysis was conducted on  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  over 25–950 °C (Fig. 4). Similarly, when measured at low temperature (<750 °C), only the cubic



**Fig. 7.** The SEM images of  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramics sintered at different temperatures.

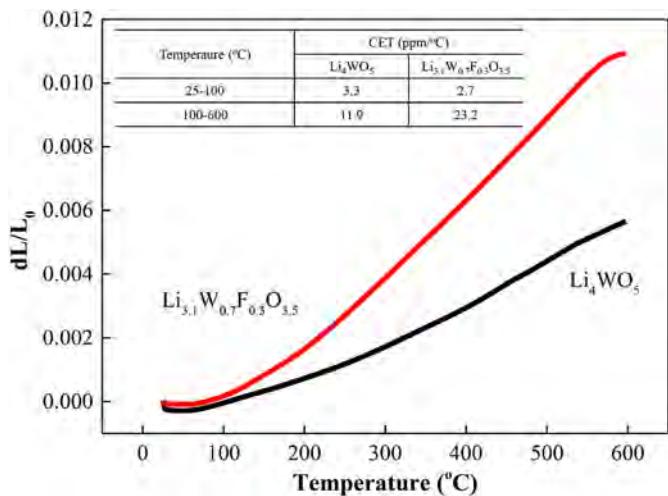


**Fig. 8.** The variations of the relative densities and microwave dielectric properties ( $\epsilon_r$ ,  $Q \times f$ , and  $\tau_f$ ) of  $(1-x)\text{Li}_4\text{WO}_5-x\text{LiF}$  ( $x = 0.1, 0.2, 0.3, 0.4, 0.5$ ) ceramics sintered at different temperatures. And the illustration of Fig. 8(a–c) gives the best performance as a function of LiF concentration.

**Table 1**

Comparison of microwave dielectric properties of some compounds with rock salt structure sintered at the densification temperature.

Composition	S.T. (°C)	$\epsilon_r$	$Q \times f$ (GHz)	$\tau_f$ (ppm/°C)	References
$\text{Li}_4\text{WO}_5$	890	8.6	23100	-2.6	[8]
$\text{Li}_3\text{NbO}_4$	1150	16.4	47100	-45	[17]
$0.3\text{Li}_3\text{NbO}_4 + 0.7\text{Li}_2\text{SnO}_3$	1200	13.5	61600	29	[28]
$\text{Li}_2\text{ZrO}_3 + 4\text{MgO}$	1500	12.65	165924	-34.66	[29]
$\text{Li}_2\text{TiO}_3 + 4\text{MgO}$	1600	13.43	233600	-27.2	[4]
$0.7\text{Li}_4\text{WO}_5 - 0.3\text{LiF}$	930	16.6	71000	-16.4	This work

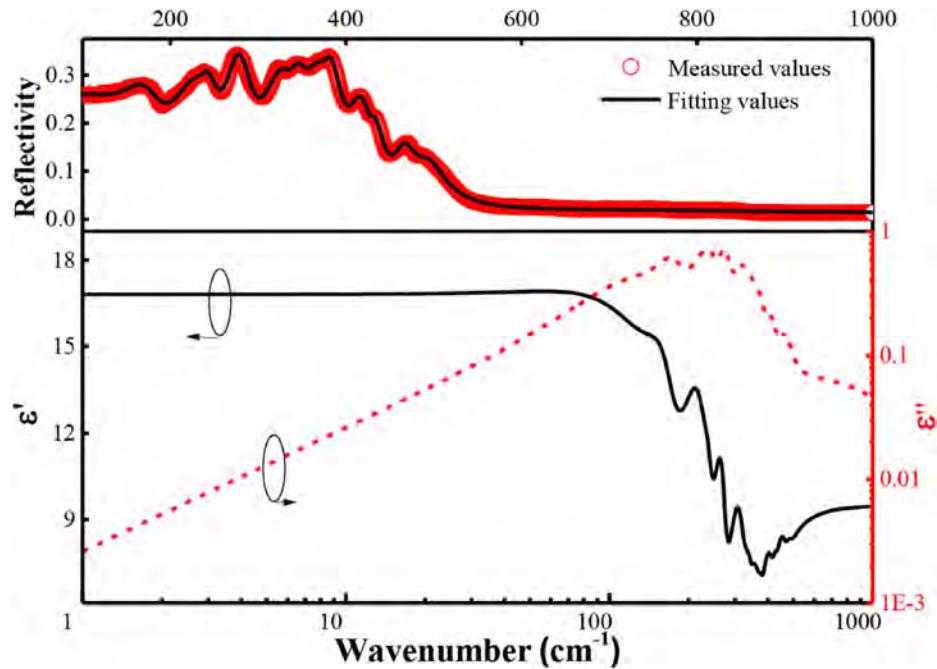


**Fig. 9.** The dilatometric curves of  $\text{Li}_4\text{WO}_5$  and  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  sintered at 890 °C and 930 °C.

$\text{Li}_4\text{WO}_5$  phase was detected while the triclinic  $\text{Li}_4\text{WO}_5$  emerged at

800 °C and became dominated at 920 °C, finally a single triclinic  $\text{Li}_4\text{WO}_5$  phase was obtained at 930 °C. When cooling down from elevated temperatures, the triclinic phase was stabilized up to room temperature, which was consistent with the irreversible feather of the phase transition. Additionally, obvious low-angle shift in XRD peaks was observed during heating process, while high-angle shift when cooling down, suggesting  $\text{Li}_4\text{WO}_5$ -based ceramics are materials with positive thermal expansion. Meanwhile, Fig. 5(a) and (b) show Rietveld refinement patterns and schematics of crystal structure for the cubic and triclinic phase of  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$ , respectively. Structural models with a space group Fm-3m and P-1 were used for refinement, respectively. The results of the Rietveld refinement using the powder XRD data are listed in Table S1 (Supplementary information). Fig. 5(c) shows the cubic phase of  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  from the selected area electron diffraction (SAED) images. The distances of the diffraction spots measured are 2.06 Å and 1.45 Å corresponding to  $d_{400}$  and  $d_{440}$  spacing, respectively. Fig. 5(d) shows the triclinic phase of  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$ , the distances of the diffraction spots are 2.20 Å and 2.19 Å corresponding to  $d_{\bar{1}12}$  and  $d_{21\bar{1}}$  spacing, respectively.

In order to determine whether LiF existed and the phase transition process of  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  occurred. Raman spectra was conducted on  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramic, as shown in Fig. 6. The



**Fig. 10.** The infrared reflectivity spectra of the  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramics in the range  $100 \text{ cm}^{-1}$ – $1000 \text{ cm}^{-1}$ .

phase transition temperature from cubic (the stretch mode marked as peach shape) to the triclinic (the stretch mode marked as circular point) phase is between  $915^\circ\text{C}$  and  $930^\circ\text{C}$ . No Raman mode of LiF exists in the temperature range of  $670^\circ\text{C}$ – $960^\circ\text{C}$ . At the same time,  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  phase transition occurs with the increase of sintering temperature. The half-peak width of the whole stretching vibration peak of Raman becomes narrower, which indicates that the order degree of high temperature phase is higher than that of low temperature phase [21]. The effect of LiF content on the transformation temperature of  $(1-x)\text{Li}_4\text{WO}_5-x\text{LiF}$  ( $x = 0.1, 0.2, 0.4, 0.6, 1.0$ ) is explained as follow. As the content of LiF increases, the order degree of oxyfluorides decreases and the transformation from cubic to triclinic phase is inhibited. Hence, higher temperature is needed to enhance the order degree up to critical point and then phase transition could occur. The similar phenomenon of rock salt  $\text{Li}_2\text{TiO}_3$ – $\text{MgO}$  ceramics has been reported [19], with the decrease of order degree by forming solid solutions. The difference is that the phase transition of  $\text{Li}_2\text{TiO}_3$ – $\text{MgO}$  binary system is from monoclinic to cubic phase. Therefore, the temperature of the order-disorder phase transition falls rapidly with the increase of MgO content.

Fig. 7 gives the SEM images of the  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramics sintered at different temperatures. Microstructural evolution is clearly distinguished with obvious grain growth with the increase of sintering temperature. Particularly, compared with the nominal  $\text{Li}_4\text{WO}_5$  composition, the microstructure of LiF-doped samples was significantly improved. Dense microstructures with well-developed grains and distinct grain boundaries were achieved. The dynamic mechanism of densification improvement is unclear yet, but it is expected that the low melting point of LiF is partially responsible.

Fig. 8 shows the variations in the relative densities and microwave dielectric properties ( $\epsilon_r$ ,  $Q \times f$ , and  $\tau_f$ ) of  $(1-x)\text{Li}_4\text{WO}_5-x\text{LiF}$  ( $x = 0.1, 0.2, 0.3, 0.4, 0.5$ ) ceramics sintered at different temperatures. The relative densities of all the compositions exhibit substantial dependence on sintering temperature. It shows a steady increase followed by a dramatic drop after approaching the highest density. Importantly, all the samples achieved relative densities

higher than 95% at their respective optimum sintering temperatures.

The variation trend of the  $\epsilon_r$  and  $Q \times f$  with sintering temperature is consistent with the density, proving that density exert prominent influences on the microwave dielectric properties (Fig. 8(b) and (c)). The compositional dependences of  $\epsilon_r$  and  $Q \times f$  values are presented in the inset of Fig. 8(b) and (c). When  $x = 0.3$ , the optimum microwave dielectric performances with  $\epsilon_r = 16.6$ , and  $Q \times f = 71,000 \text{ GHz}$  were obtained. The permittivity at microwave frequencies relies on the relative densities, secondary phases and ionic polarizability [22,23]. In our previous study, the theoretical relative permittivity value  $\sim 10.15$  of  $\text{Li}_4\text{WO}_5$  was calculated by Clausius-Mossotti equation [8]. Generally, the factors effecting the dielectric loss could be divided into two major components: the intrinsic loss and extrinsic loss relying on impurities, grain boundaries, density, lattice defects, etc [24]. In the present work, the high  $Q \times f$  value usually corresponds to the high densification. Interestingly, the  $\tau_f$  value steadily decreases with the increase of  $x$  content from  $-7.12 \text{ ppm}/^\circ\text{C}$  at  $x = 0.1$  to  $-35.11 \text{ ppm}/^\circ\text{C}$  at  $x = 0.5$ , as shown in Fig. 7(d). This phenomenon may attribute to  $\tau_f = -3 \text{ ppm}/^\circ\text{C}$  of pure  $\text{Li}_4\text{WO}_5$  and  $\tau_f = -120 \text{ ppm}/^\circ\text{C}$  of LiF. Hence, on the basis of the classical phase rule,  $\tau_f$  decreases with the increase of LiF content [25].

The microwave dielectric properties of some compounds with rock salt structure are listed in Table 1. Compared with the pure  $\text{Li}_4\text{WO}_5$ , the  $Q \times f$  value of  $\text{Li}_{2.8}\text{W}_{0.6}\text{F}_{0.4}\text{O}_3$  dramatically increases due to the distinct improvement of density. Besides,  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  has lower sintering temperature and better microwave dielectric properties in rock salt structure compounds. In summary,  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramic is a promising candidate in LTCC technology.

The dilatometric curves of  $\text{Li}_4\text{WO}_5$  and  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramics are shown in Fig. 9. Compared with the nominal  $\text{Li}_4\text{WO}_5$ ,  $\text{Li}_{3.1}\text{W}_{0.7}\text{F}_{0.3}\text{O}_{3.5}$  ceramic has a larger coefficient of thermal expansion (CTE) from the dilatometric curves. At lower temperature range ( $20^\circ\text{C}$ – $100^\circ\text{C}$ ), relatively lower thermal expansion with a CTE around  $3 \text{ ppm}/^\circ\text{C}$  were obtained, then increased to about

23 ppm/°C at elevated temperature up to 600 °C for Li<sub>3.1</sub>W<sub>0.7</sub>F<sub>0.3</sub>O<sub>3.5</sub>. This CTE value is close to that of Ag (19.5 ppm/°C), commonly used as electrode in LTCC technology, which is of significance to render good thermal expansion match between the ceramic matrix and Ag electrode in practical application [26].

Fig. 10 presents the infrared reflectivity spectra of the Li<sub>3.1</sub>W<sub>0.7</sub>F<sub>0.3</sub>O<sub>3.5</sub> ceramics in the range of 100–1000 cm<sup>-1</sup>. These spectra were analyzed and fitted using the classical harmonic oscillator model and the Fresnel formula [22,27]:

$$\epsilon^*(\omega) = \epsilon_\infty + \sum_{j=1}^n \frac{\omega_{pj}^2}{\omega_{oj}^2 - \omega^2 - ir_j\omega} \quad (2)$$

$$R(\omega) = \left| \frac{1 - \sqrt{\epsilon^*(\omega)}}{1 + \sqrt{\epsilon^*(\omega)}} \right|^2 \quad (3)$$

where  $\epsilon^*(\omega)$  is complex dielectric function, and  $\epsilon_\infty$  is the dielectric constant at the high-frequency limit;  $\omega_{oj}$ ,  $r_j$ , and  $\omega_{pj}$  are the transverse frequency, damping factor and the plasma frequency of the  $j$ -th Lorentz oscillator, respectively;  $n$  and  $R(\omega)$  are the number of transverse phonon modes and the IR reflectivity. The infrared spectra can be well fitted and phonon parameters are shown in Table S3 (Supplementary information). Moreover, the calculated permittivity  $\epsilon'$  and the dielectric loss  $\epsilon''$  obtained from the fits of the infrared reflectivity. These values match well with those measured using TE<sub>011</sub> method.

#### 4. Conclusions

(1-x)Li<sub>4</sub>WO<sub>5</sub>-xLiF ( $x = 0.1, 0.2, 0.3, 0.4, 0.5$ ) ceramics were prepared by conventional solid-state reaction method. A series of novel oxyfluorides in the Li<sub>4</sub>WO<sub>5</sub>–LiF binary system were prepared by adding an appropriate amount of LiF into Li<sub>4</sub>WO<sub>5</sub>. The temperature of phase transition (cubic-triclinic phase) moved towards high temperature with the increase of LiF content. Li<sub>3.1</sub>W<sub>0.7</sub>F<sub>0.3</sub>O<sub>3.5</sub> ceramics (96.4%) sintered at 930 °C for 4 h could be obtained with excellent microwave dielectric properties ( $Q \times f = 71,000$  GHz,  $\epsilon_r = 16.6$ ,  $\tau_f = -16.4$  ppm/°C). The infrared reflectivity spectra of the Li<sub>3.1</sub>W<sub>0.7</sub>F<sub>0.3</sub>O<sub>3.5</sub> ceramics revealed that the quality factor of Li<sub>4</sub>WO<sub>5</sub> was drastically improved by doping LiF. Therefore, the Li<sub>4</sub>WO<sub>5</sub>–LiF binary system was a promising candidate for microwave electronic device applications.

#### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

#### CRediT authorship contribution statement

**Kai Xiao:** Data curation, Writing - original draft. **Chunchun Li:** Conceptualization, Methodology. **Ying Tang:** Supervision, Formal analysis, Writing - review & editing. **Yunfei Tian:** Data curation. **Changzhi Yin:** Data curation. **Junqi Chen:** Data curation. **Jie Li:** Formal analysis, Writing - review & editing. **Lian Duan:** Data curation. **Huaicheng Xiang:** Data curation. **Liang Fang:** Resources, Supervision, Funding acquisition.

#### Acknowledgments

This work was financially supported by the Natural Science Foundation of China (Nos. 21761008 and 21965009), the Natural

Science Foundation of Guangxi Zhuang Autonomous Region (Nos. 2018GXNSFAA138175, and 2018GXNSFBA281093), the Project of Scientific Research and Technical Exploitation Program of the Guangxi Zhuang Autonomous Region (No. AA18118008, AA18118034 and AA18118023), and Projects of Education Department of Guangxi Zhuang Autonomous Region (No. 2018KY0255) and High level innovation team and outstanding scholar program of Guangxi Institutes. The author thanks the administrator of the IR beamline workstation at the National Synchrotron Radiation Laboratory (NSRL) for their assistance in IR measurements and Dr. Hongbo Cui of the Guilin University of Technology for his help with the TEM measurements.

#### Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.jallcom.2020.155320>.

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